

Resonant generation and energetics of wind-forced near-inertial motions in a geostrophic flow

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D. B. Whitt and L. N. Thomas, 2015: Resonant Generation and Energetics of Wind-Forced Near-Inertial Motions in a Geostrophic Flow. *J. Phys. Oceanogr.*, 45, 181–208.

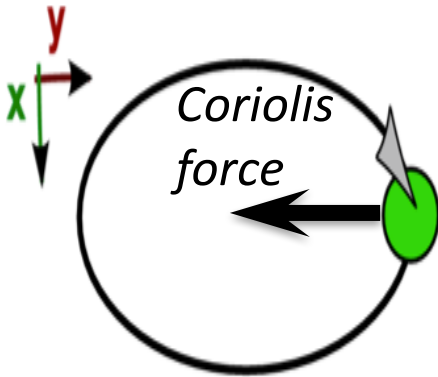
Inertial Oscillations

Inertial Frequency:

$$f = 2\Omega \sin(\text{latitude})$$

Ω = angular frequency of Earth $\sim 7.3\text{E-}5 \text{ rad/s}^{-1}$

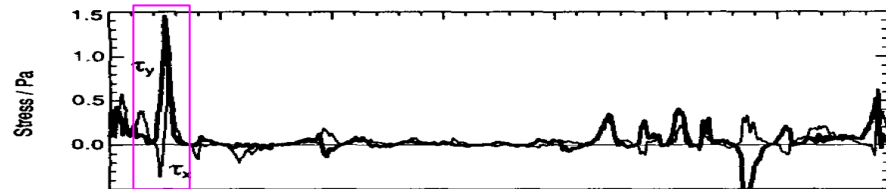
Force Diagram



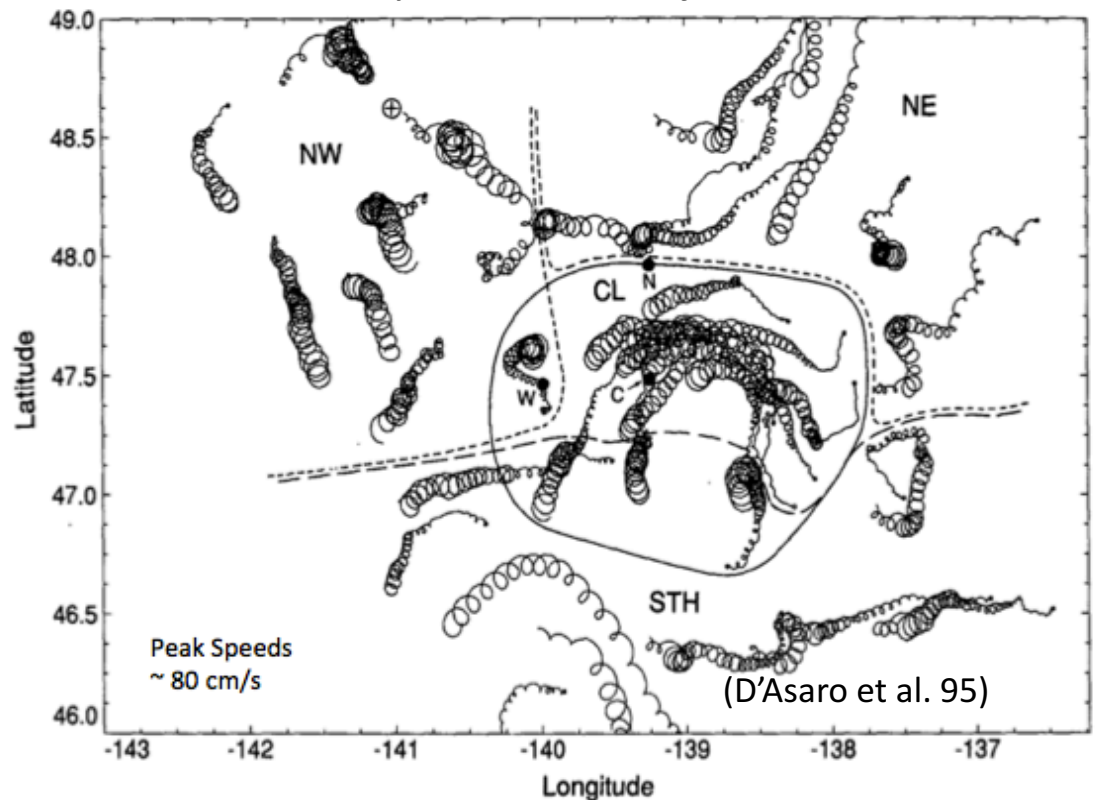
$$\frac{\partial u}{\partial t} - fv = 0$$

$$\frac{\partial v}{\partial t} + fu = 0$$

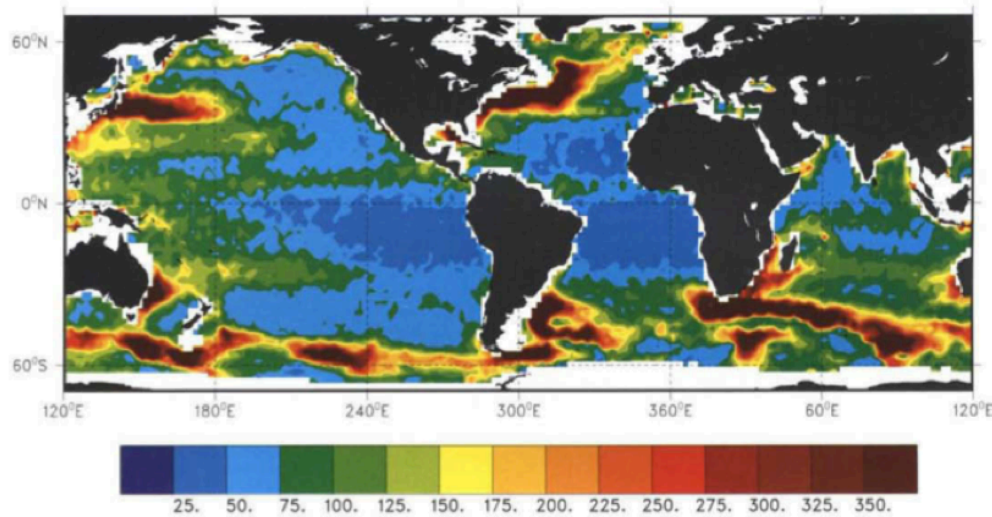
35 days of surface wind stress



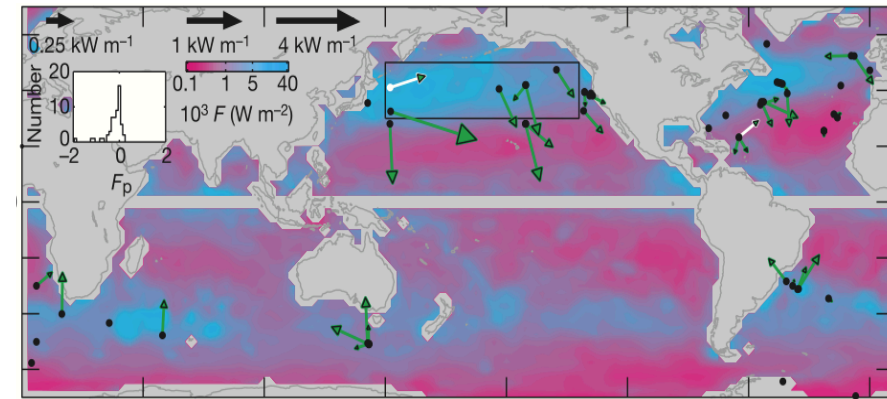
25 days of drifter trajectories



Mesoscale eddy kinetic energy (EKE) [Wunsch 2002]



Energy flux from winds to mixed-layer near-inertial motions [Alford 2003]

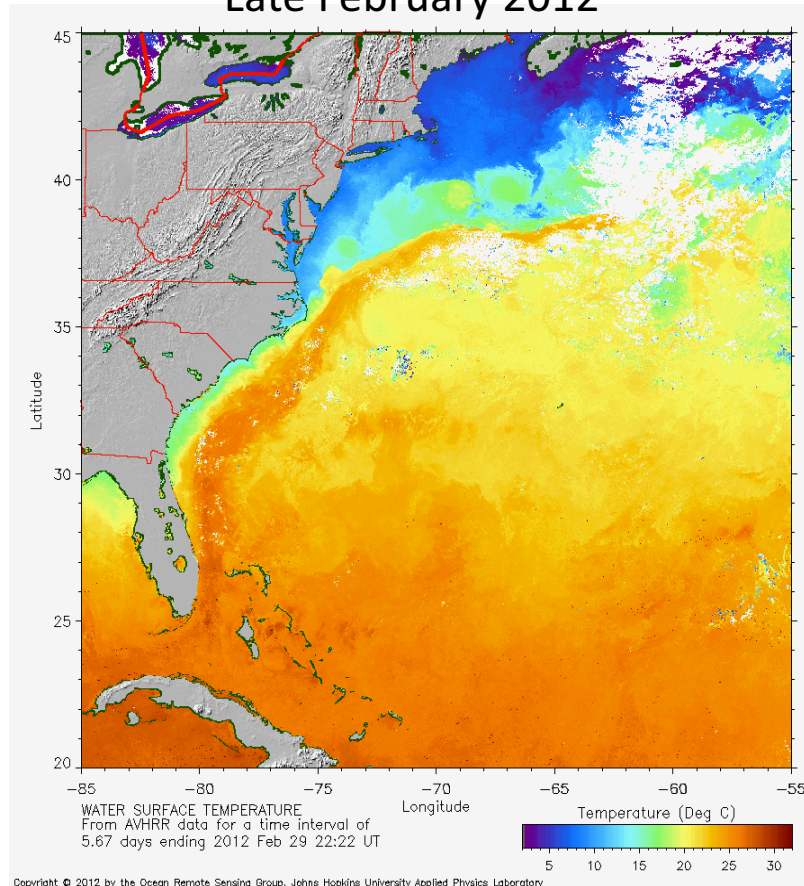


Blue is larger/pink is smaller!

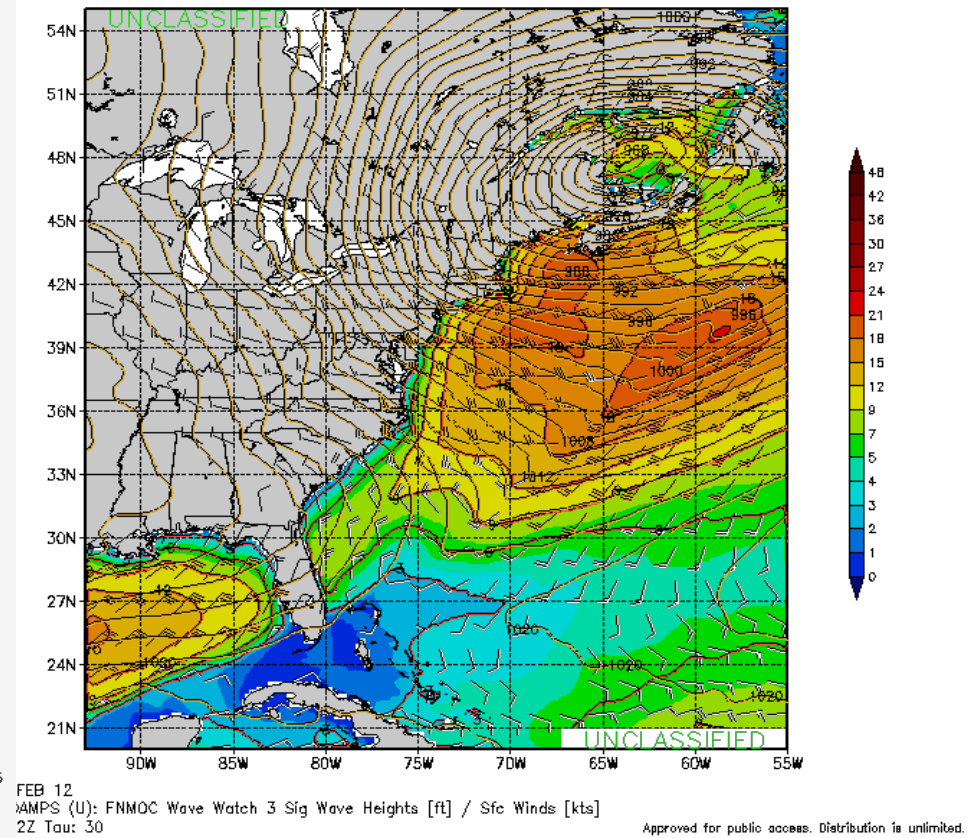
- ~ 1 TW flux into inertial oscillations.
- Much of the energy enters in regions of high EKE

Atmospheric mesoscale $\sim 100\text{-}1000\text{ km}$, Oceanic mesoscale $\sim 10\text{-}100\text{ km}$

Sea Surface Temperature,
Late February 2012



Surface Weather Chart,
Late February 2012

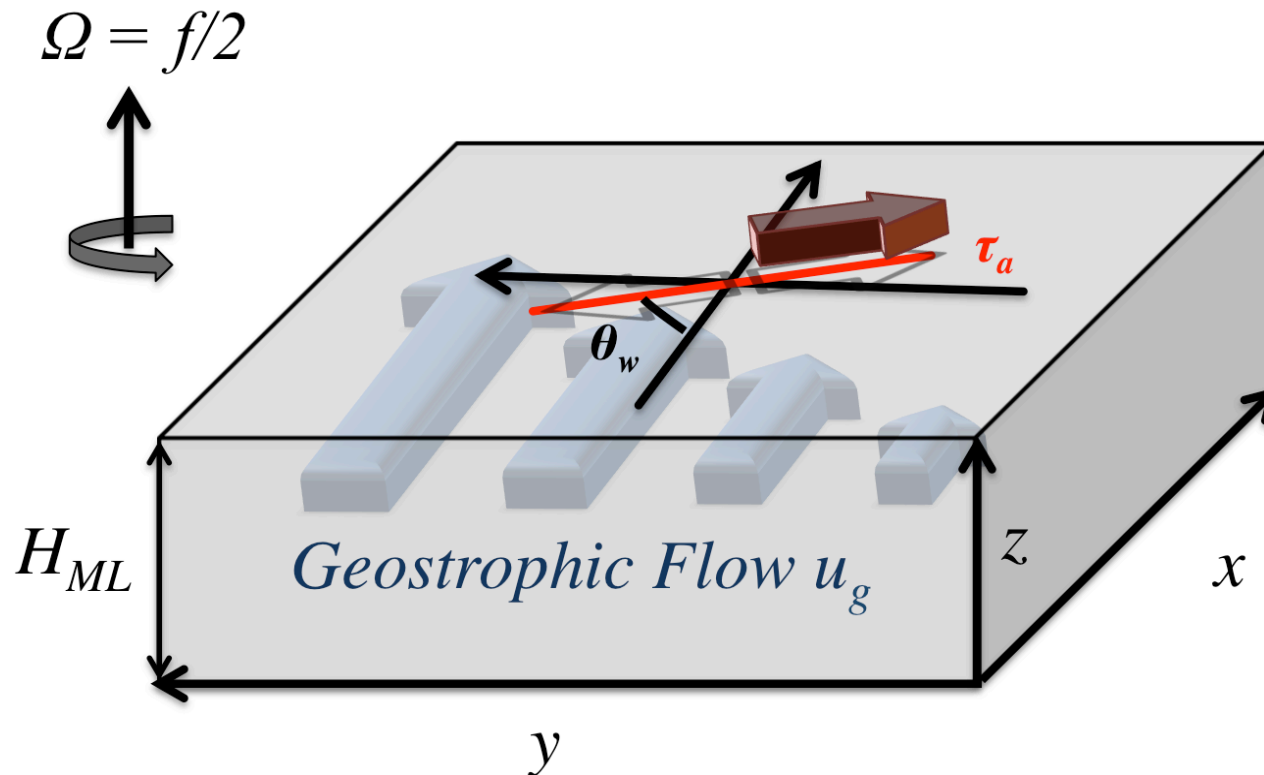


Scale Separation

How does an axisymmetric jet $u_g(y)$ modify generation of near-inertial motions by spatially-uniform oscillatory winds?

Wind stress oscillates along a line

$$\omega_{wind} = \text{constant}$$



What is an appropriate model when $du/dy \sim f$?

How does wave amplitude depend on: du/dy , θ , ω ?

Dimensionless Perturbation Momentum Eqns.

Scale separation -> model reduction

$$\frac{\partial u'}{\partial t'} + \text{Ro}_a \left(v' \frac{\partial u'}{\partial y'} + w' \frac{\partial u'}{\partial z'} \right) - v' \left(1 - \text{Ro}_g \frac{\partial u'_g}{\partial y'} \right)$$

$$+ \text{Fr}_g \text{Bu}_a^{1/2} w' \frac{\partial u'_g}{\partial z'} = \text{Ek}_a X',$$

$$\frac{\partial v'}{\partial t'} + \text{Ro}_a \left(v' \frac{\partial v'}{\partial y'} + w' \frac{\partial v'}{\partial z'} \right) + u' + \text{Bu}_a \frac{\partial p'}{\partial y'} = \text{Ek}_a Y'$$

Wave scales $\ll 1$

$$\text{Bu}_a = \left(\frac{\tilde{N} \tilde{H}_a}{f \tilde{L}_a} \right)^2 \quad \text{Ro}_a = \frac{\tilde{U}_a}{f \tilde{L}_a}$$

Larger geostrophic scale

$$\bar{\text{Ro}}_g = \tilde{U}_g / f \tilde{L}_g \sim 1$$

Modeling the mixed-layer-average momentum in a jet: a local slab model

$$\begin{aligned}
 \frac{\partial U_{ML}}{\partial t} - \frac{F^2}{f} V_{ML} &= \underbrace{-r U_{ML}}_{\substack{\text{Parameterized} \\ \text{Radiative} \\ \text{Decay/Local} \\ \text{Dissipation}}} + \underbrace{\frac{\tau_x}{\rho H_{ML}}}_{\substack{\text{Wind} \\ \text{Stress/Body} \\ \text{Force}}} \\
 \frac{\partial V_{ML}}{\partial t} + f U_{ML} &= -r V_{ML} + \frac{\tau_y}{\rho H_{ML}}
 \end{aligned}$$

$U_{ML}, V_{ML}, \tau_x, \tau_y$ *function of time only*

$$F^2 = f(f - \partial u_g / \partial y) = f^2(1 + Ro_g)$$

Classic Driven Harmonic Oscillator

$$\frac{\partial^2 U_{ML}}{\partial t^2} + 2r \frac{\partial U_{ML}}{\partial t} + (r^2 + F^2) U_{ML} = \frac{1}{\rho_0 H_{ML}} \left(\frac{\partial \tau_x}{\partial t} + r \tau_x + \frac{F^2}{f} \tau_y \right)$$

*This system is under-damped and susceptible to **resonance**.*

General Solution

$$U_{ML}(t) = U^H(t) + U^P(t)$$

Unforced Part

$$U_{ML}^H = D_1 e^{-rt} \cos(Ft) + D_2 e^{-rt} \sin(Ft)$$

Forced Part

$$U_{ML}^P(t) = A_U \sin(\omega_w t + \phi_{UP})$$

Inviscid Initial Value Problem

Resonant frequency $F = \sqrt{f \left(f - \frac{\partial u_g}{\partial y} \right)}$

Dynamics

$$\frac{\partial U_{ML}}{\partial t} - \frac{F^2}{f} V_{ML} = 0,$$

$$\frac{\partial V_{ML}}{\partial t} + f U_{ML} = 0.$$

Energetics

$$E_{ML}(t) = \frac{U_{ML}(t)^2 + V_{ML}(t)^2}{2}$$

$$= E_{ML}(0) + \underbrace{\int_0^t -\frac{\partial u_g}{\partial y} V_{ML}(s) U_{ML}(s) ds}_{\text{LSP}}$$

Lateral Shear Production



$$\frac{\partial^2 \eta}{\partial t^2} + F^2 \eta = 0$$

$$\eta(T) = \int_0^T v dt$$

*Equivalent
Lagrangian
Description*



Constant wave activity $A \neq E$

$$A = v^2 + \frac{u^2}{1 + Ro_g} = \eta_t^2 + F^2 \eta^2$$

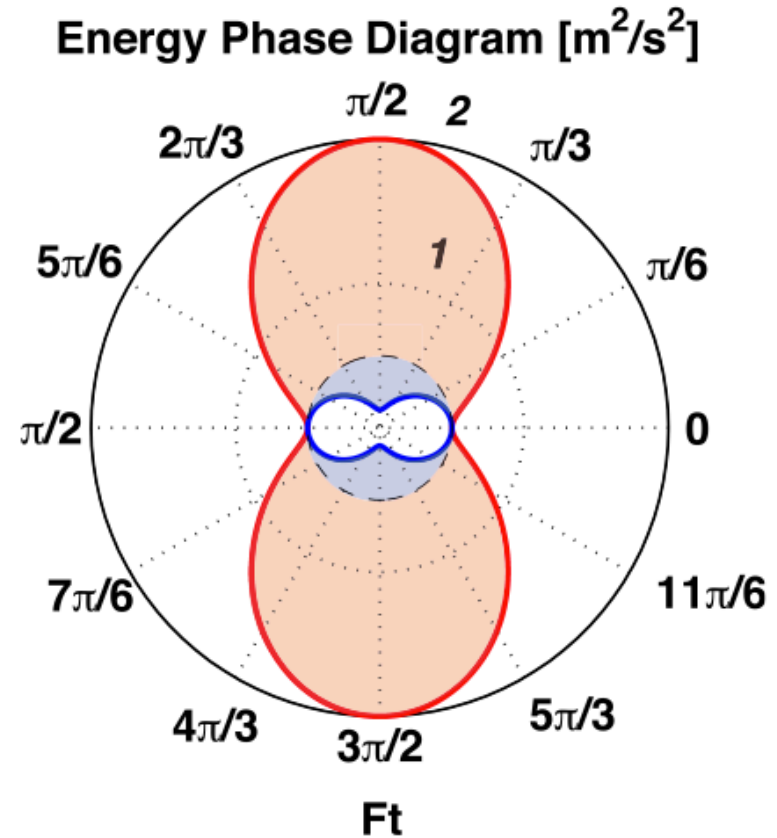
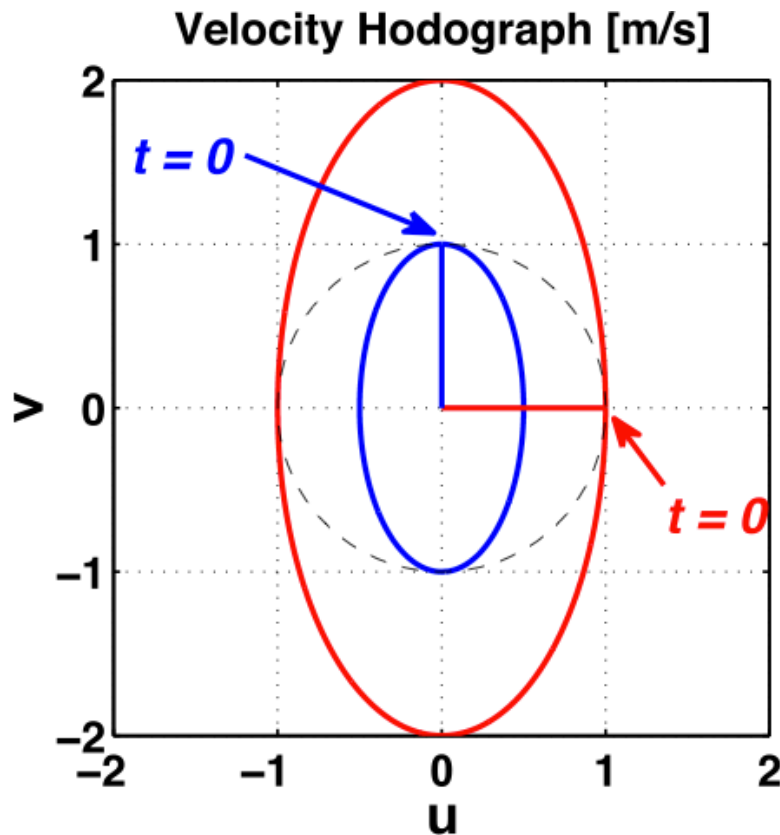
$$\langle v^2 \rangle = F^2 \langle \eta^2 \rangle \quad \text{Phase/ensemble avg.}$$

Wave activity equipartition

Example Inviscid Initial Value Problem

Resonant frequency

$$F = \sqrt{f \left(f - \frac{\partial u_g}{\partial y} \right)} = .5f$$



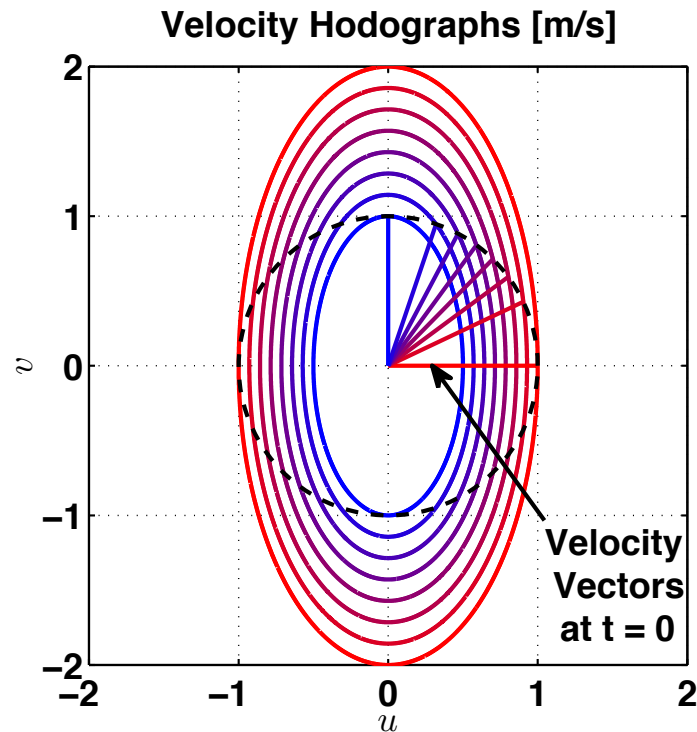
- Waves exchange energy with the background flow
- Velocity hodographs are elliptic.
- Wave energy returns to its initial value at $Ft = 2\pi$.

Ensemble of Inviscid Initial Value Problems

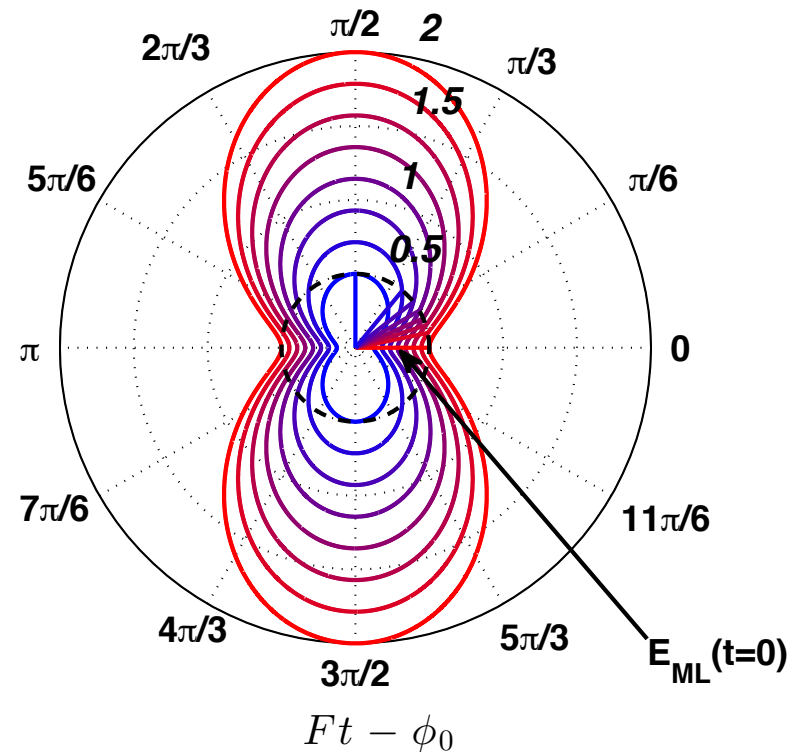
Resonant frequency

$$F = \sqrt{f \left(f - \frac{\partial u_g}{\partial y} \right)} = .5f$$

(A)



Energy Phase Diagrams [m^2/s^2]

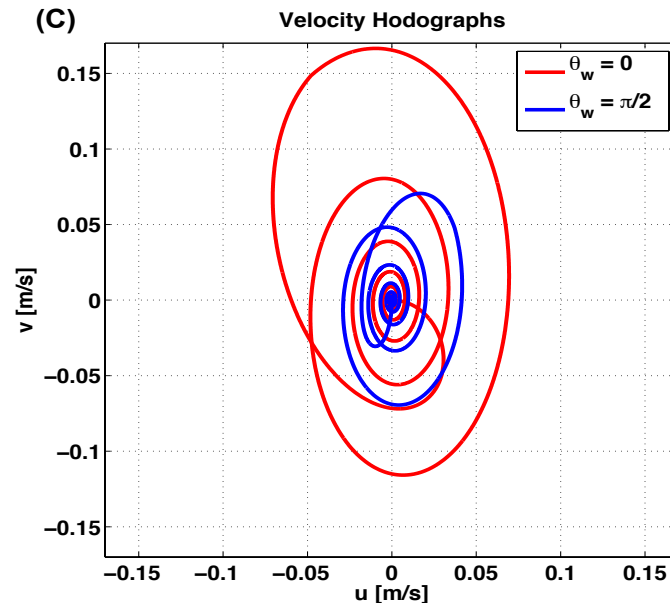
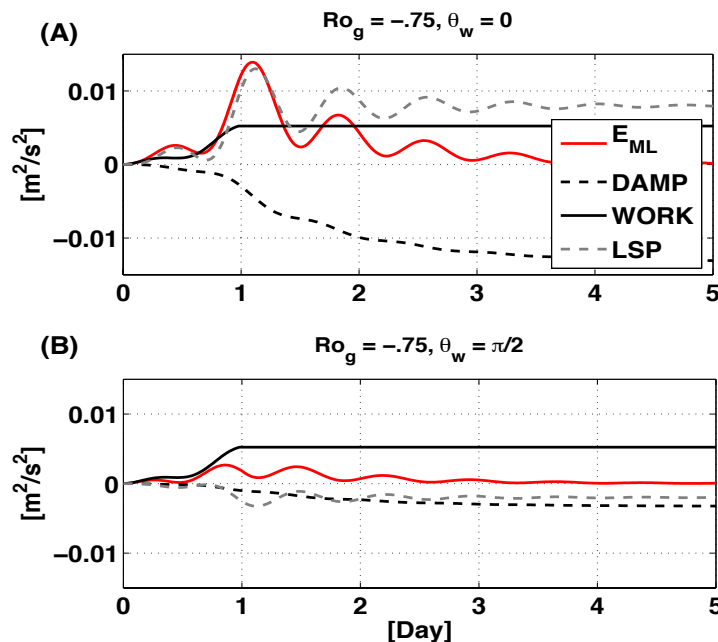


- Time and ensemble-averaged wave energy is greater than the initial energy.

Transient Forced and Damped Problem

- Forcing + dissipation drive time-integrated energy exchange between waves and geostrophic flow via LSP.
- Sign + magnitude of energy exchange depend on wind direction and geostrophic vorticity.

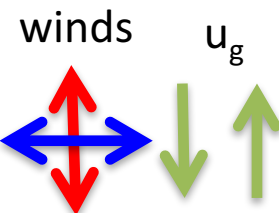
$$E_{ML}(t) - E_{ML}(0) = \underbrace{\int_0^t -\frac{\partial u_g}{\partial y} V_{ML}(s) U_{ML}(s) ds}_{\text{LSP}} + \underbrace{\int_0^t -2\tau E_{ML}(s) ds}_{\text{DAMP}} + \underbrace{\int_0^t \frac{\tau(s) \cdot \mathbf{U}_{ML}(s)}{\rho_0 H_{ML}} ds}_{\text{WORK}}$$



Resonant frequency

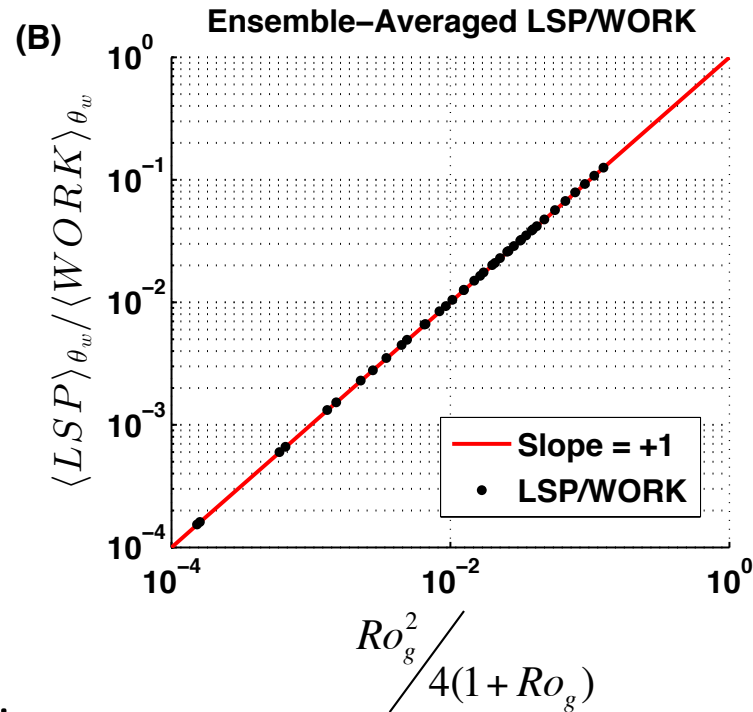
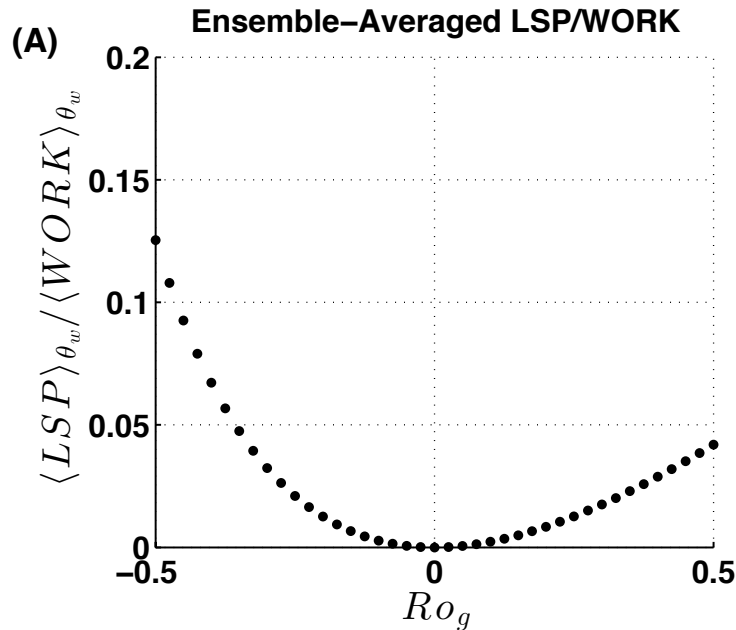
$$F = \sqrt{f \left(f - \frac{\partial u_g}{\partial y} \right)} = .5f$$

$$\omega_{\text{wind}} = F$$



Ensemble of Transient Forced Problems

forced by winds with an isotropic distribution of angles



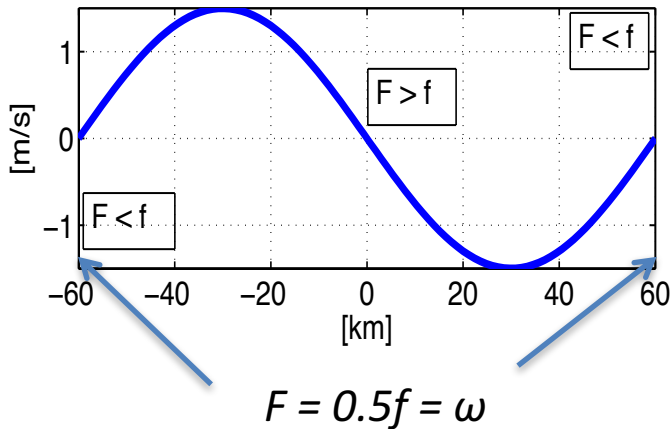
- Ensemble averaged LSP always positive
=> net energy transfer from mean to an ensemble of waves.
- Exchange is larger for anti-cyclonic vorticity
=> will tend to skew vorticity distribution of geostrophic flow toward positive values

$$\langle LSP \rangle_{\theta} \sim \langle WORK \rangle_{\theta} \frac{Ro_g^2}{4(1 + Ro_g)}$$

Some Numerical Simulations

Hydrostatic Boussinesq Equations

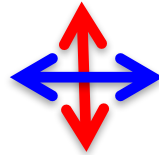
Jet Velocity



24 h Oscillatory
Wind Forcing

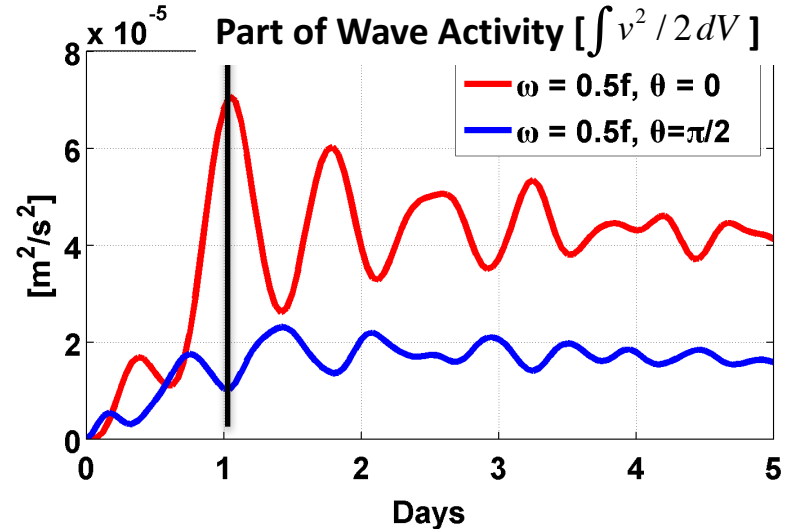
$$\theta = 0, \theta = \pi/2$$

$$\omega = 0.5f$$

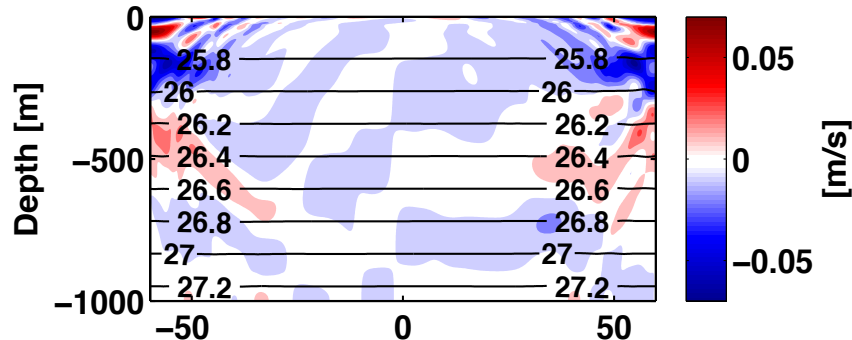


Forced

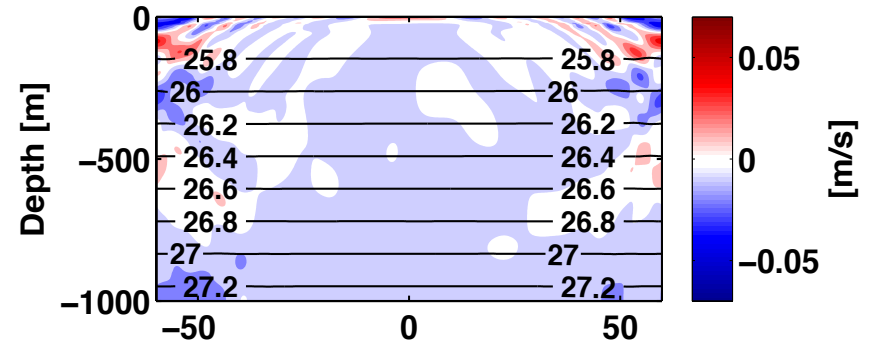
Unforced Spin-down



Cross-stream velocity, $\theta=0$, @ 5 days



Cross-stream velocity, $\theta=\pi/2$, @ 5 days



$$f = 1 \times 10^{-4} s^{-1}, N^2 = 2 \times 10^{-5} s^{-2}$$

Linearized wave activity is conserved in the absence of transience/diabatic/viscous effects

Constant time-dependent
wave activity

Initial wave activity

$$\int_V A dV = \int_V \eta_t^2 + F^2 \eta^2 + N^2 \xi^2 dV = \int_V v_0^2 + \frac{u_0^2}{1 + Ro_g} dV$$

$$\eta(T) = \int_0^T v dt, \xi(T) = \int_0^T w dt$$

$$\langle v^2 \rangle = \langle F^2 \eta^2 + N^2 \xi^2 \rangle$$

Wave activity equipartition
in the ensemble/phase average.

Easily generalizes for a baroclinic flow

(see Whitt and Thomas 2013 JPO)

Some Broader Implications

- Near-inertial waves forced by isotropic winds dissipate $O(Ro_g^2)$ fraction of energy from the mean flow at small Ro .
- At **$O(1)$ Ro**
 - KE input from winds to unbalanced motions may catalyze extraction of KE from balanced flow in proportion to wind input
 - Balanced \rightarrow unbalanced KE transfer is greater in anti-cyclonic vorticity \rightarrow may contribute to cyclone/anti-cyclone asymmetry in upper ocean